

I Radiation

① We have seen

$$\phi(\vec{r}, t) = \int d^3r' \int dt' \frac{1}{|\vec{r} - \vec{r}'|} \delta\left[t - t' - \frac{1}{c}|\vec{r} - \vec{r}'|\right] \rho(\vec{r}', t')$$

$$\vec{A}(\vec{r}, t) = \frac{1}{c} \int d^3r' \int dt' \frac{1}{|\vec{r} - \vec{r}'|} \delta\left[t - t' - \frac{1}{c}|\vec{r} - \vec{r}'|\right] \vec{J}(\vec{r}', t')$$

② Note that $\vec{r}'$ is a dummy integration variable and thus has no t or t' dependence. Only for a point particle, after we integrate over $\vec{r}'$ does $\vec{r}_p(t)$ take over the role of $\vec{r}'$. With this observation it is possible to integrate over t' and get

$$\phi(\vec{r}, t) = \int d^3r' \frac{1}{|\vec{r} - \vec{r}'|} \rho(\vec{r}', t - \frac{1}{c}|\vec{r} - \vec{r}'|)$$

$$\vec{A}(\vec{r}, t) = \frac{1}{c} \int d^3r' \frac{1}{|\vec{r} - \vec{r}'|} \vec{J}(\vec{r}', t - \frac{1}{c}|\vec{r} - \vec{r}'|)$$

- ③ Case (i) $|\vec{r} - \vec{r}'| \ll c\tau$ $\rightarrow$ a characteristic time involved in the process

In this case

$$t - \frac{|\vec{r} - \vec{r}'|}{c} \approx t$$

and thus we get

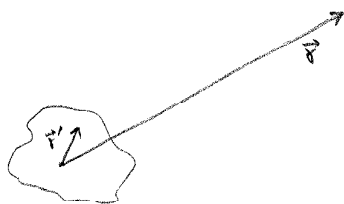
$$\phi(\vec{r}, t) = \int d^3r' \frac{1}{|\vec{r} - \vec{r}'|} \rho(\vec{r}', t)$$

$$\vec{A}(\vec{r}, t) = \frac{1}{c} \int d^3r' \frac{1}{|\vec{r} - \vec{r}'|} \vec{J}(\vec{r}', t)$$

which is very close to static.

- Case (ii) $|\vec{r}| \gg |\vec{r}'|$ (RADIATION)

Observation point $\vec{r}$ is very far from source distributed $\vec{r}'$.



$$\begin{aligned} |\vec{r} - \vec{r}'| &= \sqrt{r^2 + r'^2 - 2\vec{r} \cdot \vec{r}'} \\ &= r \sqrt{1 + \frac{r'^2}{r^2} - 2\frac{\hat{r} \cdot \vec{r}'}{r}} \\ &= r \left[1 - \frac{\hat{r} \cdot \vec{r}'}{r} + O\left(\frac{1}{r^2}\right) \right] \end{aligned}$$

④ Thus in the regime of radiation

$$|\vec{r} - \vec{r}'| = r - \hat{r} \cdot \vec{r}' + O\left(\frac{1}{r}\right)$$

$$\frac{1}{|\vec{r} - \vec{r}'|} = \frac{1}{[r - \hat{r} \cdot \vec{r}' + O(\frac{1}{r})]}$$

$$= \frac{1}{r} \frac{1}{[1 - \frac{\hat{r} \cdot \vec{r}'}{r} + O(\frac{1}{r})]}$$

$$= \frac{1}{r} \left[1 + \frac{\hat{r} \cdot \vec{r}'}{r} + O\left(\frac{1}{r}\right) \right]$$

$$= \frac{1}{r} + O\left(\frac{1}{r}\right)$$

$$t_r \equiv t - \frac{1}{c} |\vec{r} - \vec{r}'| = t - \frac{1}{c} \left[r - \hat{r} \cdot \vec{r}' + O\left(\frac{1}{r}\right) \right]$$

$$= t - \frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c} + O\left(\frac{1}{r}\right)$$

$$\textcircled{5} \quad \phi(\vec{r}, t) = \int d^3 r' \frac{1}{|\vec{r} - \vec{r}'|} \rho(\vec{r}', t - \frac{1}{c} |\vec{r} - \vec{r}'|)$$

$$= \frac{1}{r} \int d^3 r' \rho\left(\vec{r}', t - \frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}\right) + O\left(\frac{1}{r}\right)$$

$$\vec{A}(\vec{r}, t) = \frac{1}{c} \frac{1}{r} \int d^3 r' \vec{J}\left(\vec{r}', t - \frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}\right) + O\left(\frac{1}{r}\right)$$

II Radiations fields

$$\textcircled{1} \quad \phi(\vec{r}, t) = \frac{1}{8} \int d^3r' \rho(\vec{r}', t - \frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}) + O\left(\frac{1}{r^2}\right)$$

$$\vec{A}(\vec{r}, t) = \frac{1}{c} \frac{1}{r} \int d^3r' \vec{J}(\vec{r}', t - \frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}) + O\left(\frac{1}{r^2}\right)$$

$$\textcircled{2} \quad \vec{E}(\vec{r}, t) = -\vec{\nabla}\phi - \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

$$\vec{B}(\vec{r}, t) = \vec{\nabla} \times \vec{A}$$

$$t_r \equiv t - \frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}$$

$$\textcircled{3} \quad \vec{\nabla} \left[\frac{1}{8} f\left(t - \frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}\right) \right]$$

$$= \left[\vec{\nabla} \frac{1}{8} \right] f\left(t - \frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}\right) + \frac{1}{8} [\vec{\nabla} t_r] \left\{ \frac{\partial f(t)}{\partial t} \right\}_{t=t_r}$$

$$= -\frac{\hat{r}}{8} f(t_r) + \frac{1}{8} \left\{ \frac{\partial f(t)}{\partial t} \right\}_{t=t_r} \left[-\frac{1}{c} \vec{\nabla} r + \frac{1}{c} \vec{\nabla}(\hat{r} \cdot \vec{r}') \right]$$

$$\textcircled{4} \quad \vec{\nabla}(\hat{r} \cdot \vec{r}') = \frac{1}{8} \left[\vec{r}' - \hat{r}(\hat{r} \cdot \vec{r}') \right]$$

$$= -\frac{1}{8} \hat{r} \times (\hat{r} \times \vec{r}')$$

$$\textcircled{5} \quad \text{Using } \textcircled{4} \text{ in } \textcircled{3}$$

$$\vec{\nabla} \left[\frac{1}{8} f\left(t - \frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}\right) \right] = -\frac{1}{c} \frac{\hat{r}}{8} \left\{ \frac{\partial f(t)}{\partial t} \right\}_{t=t_r} + O\left(\frac{1}{r^2}\right)$$

⑥ Using ⑤ in ②

$$\vec{B}(\vec{r}, t) = -\frac{1}{c^2} \frac{\hat{r}}{r} \times \int d^3r' \left\{ \frac{\partial \vec{J}(\vec{r}', t)}{\partial t} \right\}_{t=t-\frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}} + O\left(\frac{1}{r^2}\right)$$

$$\vec{E}(\vec{r}, t) = \frac{1}{c} \frac{\hat{r}}{r} \int d^3r' \left\{ \frac{\partial P(\vec{r}', t)}{\partial t} \right\}_{t=t_r} - \frac{1}{c^2} \frac{1}{r} \int d^3r' \underbrace{\frac{\partial t_r}{\partial t}}_{=1} \left\{ \frac{\partial \vec{J}(\vec{r}', t)}{\partial t} \right\}_{t=t_r} + O\left(\frac{1}{r^2}\right)$$

⑦ $\frac{\partial P(\vec{r}, t)}{\partial t} + \vec{\nabla} \cdot \vec{J}(\vec{r}, t) = 0$ ← usual local ^{charge} conservation condition.

$$\left\{ \frac{\partial P(\vec{r}, t)}{\partial t} \right\}_{t=t_r} = - \left\{ \vec{\nabla}_{\vec{r}'} \cdot \vec{J}(\vec{r}', t) \right\}_{t=t-\frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}}$$

↙ acts on explicit $\vec{r}'$
and on $\vec{r}'$ dependence in t

$$= - \left\{ \vec{\nabla}_{\vec{r}'} \cdot \vec{J}(\vec{r}', t) \right\}_{t=t_r} - [\vec{\nabla}_{\vec{r}'} t_r] \cdot \left\{ \frac{\partial \vec{J}(\vec{r}', t)}{\partial t} \right\}_{t=t_r}$$

↙ acts only on explicit $\vec{r}'$
and not on $\vec{r}'$ dep. in t .

$$= - \left\{ \vec{\nabla}_{\vec{r}'} \cdot \vec{J}(\vec{r}', t) \right\}_{t=t_r} + \frac{\hat{r}}{c} \cdot \left\{ \frac{\partial \vec{J}(\vec{r}', t)}{\partial t} \right\}_{t=t_r} + O\left(\frac{1}{r}\right)$$

⑧ $\int d^3r' \left\{ \frac{\partial P(\vec{r}', t)}{\partial t} \right\}_{t=t_r} = - \int d^3r' \left\{ \vec{\nabla}_{\vec{r}'} \cdot \vec{J}(\vec{r}', t) \right\}_{t=t_r} + \frac{\hat{r}}{c} \cdot \int d^3r' \left\{ \frac{\partial \vec{J}(\vec{r}', t)}{\partial t} \right\}_{t=t_r} + O\left(\frac{1}{r}\right)$

↘ = 0 ($\because \vec{J}$ is confined to finite volume)

$$= + \frac{\hat{r}}{c} \cdot \int d^3r' \left\{ \frac{\partial \vec{J}(\vec{r}', t)}{\partial t} \right\}_{t=t_r} + O\left(\frac{1}{r}\right)$$

⑨ Using ⑧ in ⑥

$$\begin{aligned}
 \vec{E}(\vec{r}, t) &= \frac{1}{c} \frac{\hat{r}}{r} \int d^3r' \left\{ \frac{\partial}{\partial t} \rho(\vec{r}', t) \right\}_{t=t_1} - \frac{1}{c^2} \frac{1}{r} \int d^3r' \left\{ \frac{\partial}{\partial t} \vec{J}(\vec{r}', t) \right\}_{t=t_1} + O\left(\frac{1}{r^2}\right) \\
 &= + \frac{1}{c^2} \frac{\hat{r}}{r} \int d^3r' \left\{ \frac{\partial}{\partial t} \vec{J}(\vec{r}', t) \right\}_{t=t_1} - \frac{1}{c^2} \frac{1}{r} \int d^3r' \left\{ \frac{\partial}{\partial t} \vec{J}(\vec{r}', t) \right\}_{t=t_1} + O\left(\frac{1}{r^2}\right) \\
 &= + \frac{1}{c^2} \frac{1}{r} \int d^3r' \left[\hat{r} \cdot \left\{ \frac{\partial}{\partial t} \vec{J}(\vec{r}', t) \right\}_{t=t_1} - \left\{ \frac{\partial}{\partial t} \vec{J}(\vec{r}', t) \right\}_{t=t_1} \cdot \hat{r} \right] + O\left(\frac{1}{r^2}\right) \\
 &= \frac{1}{c^2} \frac{1}{r} \int d^3r' \left(\hat{r} \times \left(\hat{r} \times \left\{ \frac{\partial \vec{J}(\vec{r}', t)}{\partial t} \right\}_{t=t_1} \right) \right) + O\left(\frac{1}{r^2}\right) \\
 &\quad \text{(using } \hat{r} \times (\hat{r} \times \vec{v}) = \hat{r}(\hat{r} \cdot \vec{v}) - \vec{v}) \\
 &= - \hat{r} \times \left[- \frac{1}{c^2} \frac{\hat{r}}{r} \times \int d^3r' \left\{ \frac{\partial \vec{J}(\vec{r}', t)}{\partial t} \right\}_{t=t_1} \right] + O\left(\frac{1}{r^2}\right) \\
 &= - \hat{r} \times \vec{B}(\vec{r}, t) + O\left(\frac{1}{r^2}\right) \quad \text{(using ⑥)}
 \end{aligned}$$

⑩ Thus for radiation fields we have

$$\vec{E} = - \hat{r} \times \vec{B} + O\left(\frac{1}{r^2}\right) \quad \& \quad \vec{B} = \hat{r} \times \vec{E} + O\left(\frac{1}{r^2}\right)$$

$$\text{(ii)} \quad E^2 = B^2 + O\left(\frac{1}{r^2}\right)$$

III. Radiation field due to a point charge

① We saw

$$\vec{B}(\vec{r}, t) = -\frac{1}{c^2} \hat{r} \times \int d^3r' \left\{ \frac{\partial}{\partial t} \vec{J}(\vec{r}', t) \right\}_{t=t-\frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}} + O\left(\frac{1}{r}\right)$$

$$\vec{E}(\vec{r}, t) = -\hat{r} \times \vec{B}(\vec{r}, t) + O\left(\frac{1}{r}\right)$$

② For a point charge

$$\rho(\vec{r}, t) = q \delta^{(3)}(\vec{r} - \vec{r}_q(t))$$

$$\vec{J}(\vec{r}, t) = q \vec{v}_q(t) \delta^{(3)}(\vec{r} - \vec{r}_q(t))$$

$$\textcircled{3} \quad \vec{B}(\vec{r}, t) = -\frac{q}{c^2} \hat{r} \times \int d^3r' \left\{ \frac{\partial}{\partial t} \left[\vec{v}_q(t) \delta^{(3)}(\vec{r}' - \vec{r}_q(t)) \right] \right\}_{t=t-\frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}} + O\left(\frac{1}{r}\right)$$

CAUTION: $\frac{\partial}{\partial t}$ does not come out of the r' integral because it is evaluated at a point where it has r' dependence.

⑦ Using ⑥ in ⑤

$$\int d^3r' \left\{ \frac{\partial}{\partial t} \left[\vec{V}_q(t) \delta^{(3)}(\vec{r}' - \vec{r}_q(t)) \right] \right\}_{t=t-\frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}}$$

$$= \left[\vec{a}_q(t) + \vec{a}_q(t) \left(\hat{r} \cdot \frac{\vec{V}_q(t)}{c} \right) + \frac{\vec{V}_q(t)}{c} \left(\hat{r} \cdot \vec{a}_q(t) \right) \right]_{t=t-\frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}}$$

⑧ Using ⑦ in ③

$$\vec{B}(\vec{r}, t) = -\frac{q}{c^2} \frac{\hat{r}}{r} \times \left[\vec{a}_q(t) + \vec{a}_q(t) \left(\hat{r} \cdot \frac{\vec{V}_q(t)}{c} \right) + \frac{\vec{V}_q(t)}{c} \left(\hat{r} \cdot \vec{a}_q(t) \right) \right]_{t=t-\frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}} + O\left(\frac{1}{r^2}\right)$$

$$\vec{E}(\vec{r}, t) = -\hat{r} \times \vec{B}(\vec{r}, t) + O\left(\frac{1}{r^2}\right)$$

⑨ For a non-relativistic particle ($\frac{v}{c} \ll 1$)

$$\vec{B}(\vec{r}, t) = -\frac{q}{c^2} \frac{\hat{r}}{r} \times \vec{a}_q(t) \Big|_{t=t-\frac{r}{c} + \frac{\hat{r} \cdot \vec{r}'}{c}} + O\left(\frac{1}{r^2}\right)$$

$$\vec{E}(\vec{r}, t) = -\hat{r} \times \vec{B}(\vec{r}, t) + O\left(\frac{1}{r^2}\right)$$

IV Radiation field due to a non-relativistic point particle

① We saw

$$\vec{B}(\vec{r}, t) = -\frac{q}{c^2} \frac{\hat{r}}{r} \times \int d^3r' \left\{ \frac{\partial}{\partial t} [\vec{V}_q(t) \delta^{(3)}(\vec{r}' - \vec{r}_q(t))] \right\}_{t=t-\frac{r}{c}-\frac{\hat{r} \cdot \vec{r}'}{c}} + O\left(\frac{1}{r^2}\right)$$

$$\vec{E}(\vec{r}, t) = -\hat{r} \times \vec{B}(\vec{r}, t)$$

② For $\frac{v}{c} \ll 1$

$$\begin{aligned} t_r &= t - \frac{r}{c} - \frac{\hat{r} \cdot \vec{r}'}{c} \\ &= t - \frac{r}{c} - \hat{r} \cdot \frac{\vec{V}'}{c} \tau \rightarrow \text{a time involved in the process} \\ &\approx t - \frac{r}{c} \quad (\text{for } \frac{v}{c} \ll 1) \end{aligned}$$

$$\begin{aligned} \textcircled{3} \quad \vec{B}(\vec{r}, t) &= -\frac{q}{c^2} \frac{\hat{r}}{r} \times \int d^3r' \left\{ \frac{\partial}{\partial t} [\vec{V}_q(t) \delta^{(3)}(\vec{r}' - \vec{r}_q(t))] \right\}_{t=t-\frac{r}{c}} + O\left(\frac{1}{r^2}, \frac{v}{c}\right) \\ &= -\frac{q}{c^2} \frac{\hat{r}}{r} \times \left\{ \frac{\partial \vec{V}_q(t)}{\partial t} \int d^3r' \delta^{(3)}(\vec{r}' - \vec{r}_q(t)) \right\}_{t=t-\frac{r}{c}} + O\left(\frac{1}{r^2}, \frac{v}{c}\right) \\ &= -\frac{q}{c^2} \frac{\hat{r}}{r} \times \vec{a}_q(t) \Big|_{t=t-\frac{r}{c}} + O\left(\frac{1}{r^2}, \frac{v}{c}\right) \end{aligned}$$

$$\vec{E}(\vec{r}, t) = -\hat{r} \times \vec{B}(\vec{r}, t) + O\left(\frac{1}{r^2}, \frac{v}{c}\right)$$