

Conservation of energy

- ① Let us remove the constraint of electrostatics and consider the complete Maxwell's equations

$$\vec{\nabla} \cdot \vec{D} = \rho_e$$

$$-\vec{\nabla} \times \vec{E} - \frac{\partial \vec{B}}{\partial t} = 0 \rightarrow \vec{j}_m$$

$$\vec{\nabla} \cdot \vec{B} = 0 \rightarrow \rho_m$$

$$\vec{\nabla} \times \vec{H} - \frac{\partial \vec{D}}{\partial t} = \vec{j}_e$$

$$\vec{D} = \epsilon_0 \vec{E}$$

$$\vec{B} = \mu_0 \vec{H}$$

$\rho_m \rightarrow$ magnetic monopole charge density

$\vec{j}_m \rightarrow$ magnetic monopole current density

- ② The Lorentz force is given by

$$\vec{F} = q_e [\vec{E} + \vec{v} \times \vec{B}] + q_m [\vec{H} - \vec{v} \times \vec{D}]$$

- ③ where the negative sign is suggested by the invariance under

$$\rho_e \rightarrow \rho_m$$

$$\vec{j}_e \rightarrow \vec{j}_m$$

$$\vec{E} \rightarrow \vec{H}$$

and

$$\rho_m \rightarrow -\rho_e$$

$$\vec{j}_m \rightarrow -\vec{j}_e$$

$$\vec{H} \rightarrow -\vec{E}$$

- ④ Note that

$$\text{Power} = \frac{\text{Energy}}{\text{time}} = \frac{\vec{F} \cdot d}{t} \rightarrow \vec{F} \cdot \vec{v}$$

⑤ Thus, the rate of energy transferred from the electromagnetic field to the charge is given by

$$\begin{aligned}\vec{F} \cdot \vec{v} &= q_e [\vec{E} + \vec{v} \times \vec{B}] \cdot \vec{v} + q_m [\vec{H} - \vec{v} \times \vec{D}] \cdot \vec{v} \\ &= (q_e \vec{v}) \cdot \vec{E} + (q_m \vec{v}) \cdot \vec{H} \\ &= \int d^3x [\vec{j}_e \cdot \vec{E} + \vec{j}_m \cdot \vec{H}]\end{aligned}$$

where we used the charge density of a particle as

$$\begin{aligned}\vec{j}_e(\vec{r}) &= q_e \vec{v}_a \delta^{(3)}(\vec{r} - \vec{r}_a(t)) \\ \vec{j}_m(\vec{r}) &= q_m \vec{v}_a \delta^{(3)}(\vec{r} - \vec{r}_a(t)).\end{aligned}$$

⑥ Using Maxwell's equation we have.

$$\begin{aligned}\vec{j}_e \cdot \vec{E} &= \left[\vec{\nabla} \times \vec{H} - \frac{\partial \vec{D}}{\partial t} \right] \cdot \vec{E} \\ &= (\vec{\nabla} \times \vec{H}) \cdot \vec{E} - \frac{\partial}{\partial t} \left(\frac{1}{2} \epsilon_0 E^2 \right)\end{aligned}$$

$$\begin{aligned}\vec{j}_m \cdot \vec{H} &= \left[-\vec{\nabla} \times \vec{E} - \frac{\partial \vec{B}}{\partial t} \right] \cdot \vec{H} \\ &= -(\vec{\nabla} \times \vec{E}) \cdot \vec{H} - \frac{\partial}{\partial t} \left(\frac{1}{2} \mu_0 H^2 \right)\end{aligned}$$

⑦ Adding we have

$$\vec{j}_e \cdot \vec{E} + \vec{j}_m \cdot \vec{H} + \frac{\partial}{\partial t} \left(\frac{1}{2} \epsilon_0 E^2 + \frac{1}{2} \mu_0 H^2 \right) = (\vec{\nabla} \times \vec{H}) \cdot \vec{E} - (\vec{\nabla} \times \vec{E}) \cdot \vec{H}.$$

$$\begin{aligned}
 \textcircled{8} \quad \vec{\nabla} \cdot (\vec{E} \times \vec{H}) &= \epsilon_{ijk} \nabla_i (E_j H_k) \\
 &= \epsilon_{ijk} (\nabla_i E_j) H_k + \epsilon_{ijk} E_j (\nabla_i H_k) \\
 &= (\vec{\nabla} \times \vec{E}) \cdot \vec{H} - (\vec{\nabla} \times \vec{H}) \cdot \vec{E}
 \end{aligned}$$

Using $\textcircled{8}$ in $\textcircled{7}$ we have

$$\textcircled{9} \quad \frac{\partial U}{\partial t} + \vec{\nabla} \cdot \vec{S} + \vec{j}_e \cdot \vec{E} + \vec{j}_m \cdot \vec{H} = 0$$

where

$$U = \frac{1}{2} \epsilon_0 E^2 + \frac{1}{2} \mu_0 H^2 \quad \rightarrow \text{electromagnetic field energy density}$$

$$\vec{S} = \vec{E} \times \vec{H} \quad \rightarrow \text{flux of electromagnetic energy (Poynting vector)}$$

Integrating over a volume V we have.

$$\textcircled{10} \quad \frac{\partial}{\partial t} \left(\int_V d^3r U \right) + \oint_S d\vec{a} \cdot \vec{S} + \int_V d^3r \vec{j}_e \cdot \vec{E} + \int_V d^3r \vec{j}_m \cdot \vec{H} = 0$$

rate of change of electromagnetic field energy inside volume V .

rate at which electromagnetic field energy is leaving volume V .

rate of transfer of electromagnetic field energy to the charges.

⑪ What is the interpretation in electrostatics?

$$\vec{\nabla} \cdot \vec{E} = 0$$

$$U = \frac{1}{2} \epsilon_0 E^2$$

$$W = q_e \phi$$

such that we have.

$$\frac{1}{2} \epsilon_0 E^2 + q_e \phi = \text{constant}$$

and in integrated form we have.

$$\int_V d^3r \left(\frac{1}{2} \epsilon_0 E^2 \right) + q_a \phi(\vec{r}_a) = \text{const}$$

↙
electromagnetic field
energy inside volume V .

↘
work done on the charges
to bring them from infinity.

Energy in electrostatics

① $\vec{\nabla} \cdot \vec{E} = \frac{1}{\epsilon_0} \rho$

$\vec{\nabla} \times \vec{E} = 0 \Rightarrow \vec{E} = -\vec{\nabla} \phi$

$\phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int d^3r' \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|}$

$\phi \rightarrow$ Electric potential

$\vec{E} \rightarrow$ Electric field

② Bring a test charge "q"

$\vec{F} = q_a \vec{E}(\vec{r}_a)$

$U = q_a \phi(\vec{r}_a)$

$\vec{F} \rightarrow$ force on charge q

$U \rightarrow$ interaction energy between charge q and charge distribution $\rho(\vec{r})$.

③ Relation between force and energy

$\vec{F} = q_a \vec{E}(\vec{r}_a)$

$= -q_a \vec{\nabla} \phi(\vec{r}_a)$

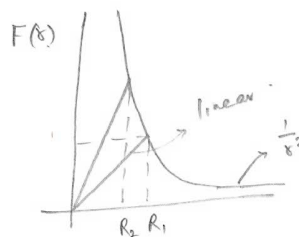
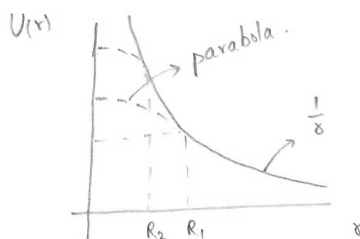
$= -\vec{\nabla} [q_a \phi(\vec{r}_a)]$

$= -\vec{\nabla} U$

④ Let us analyse the Rutherford scattering experiment.

$$V(r) = q \phi(r) = \begin{cases} \frac{qQ}{4\pi\epsilon_0} \frac{1}{r} & r > R \\ \frac{qQ}{4\pi\epsilon_0} \frac{1}{R} \left[\frac{3}{2} - \frac{1}{2} \frac{r^2}{R^2} \right] & r < R \end{cases}$$

$$\vec{F}(r) = q \vec{E}(r) = \begin{cases} \frac{qQ}{4\pi\epsilon_0} \frac{1}{r^2} \hat{r} & r > R \\ \frac{qQ}{4\pi\epsilon_0} \frac{1}{R^2} \frac{r}{R} \hat{r} & r < R \end{cases}$$



Kinetic energy of a projectile can probe the radius of a charged sphere.

— if projectile rips through:

$$\frac{qQ}{4\pi\epsilon_0} \frac{1}{R} \frac{3}{2} < \frac{1}{2} m v^2$$

$$\Rightarrow \left(\frac{qQ}{4\pi\epsilon_0} \frac{3}{2} \right) \left(\frac{1}{2} m v^2 \right)^{-1} < R$$

— if projectile bounces back:

$$\frac{qQ}{4\pi\epsilon_0} \frac{1}{R} \frac{3}{2} > \frac{1}{2} m v^2$$

$$\Rightarrow \left(\frac{qQ}{4\pi\epsilon_0} \frac{3}{2} \right) \left(\frac{1}{2} m v^2 \right)^{-1} > R$$

⑤ Point particle

$$\vec{F} = q_a \vec{E}(\vec{r}_a)$$

$$= -\vec{\nabla} U$$

$$U = q_a \phi(\vec{r}_a)$$

⑥ Electric dipole moment:

$$\vec{F} = q \vec{E}(\vec{r}_+) - q \vec{E}(\vec{r}_-)$$

$$= q \vec{E}\left(\vec{r}_0 + \frac{\vec{a}}{2}\right) - q \vec{E}\left(\vec{r}_0 - \frac{\vec{a}}{2}\right)$$

$$= q \vec{E}(\vec{r}_0) + q \frac{\vec{a}}{2} \cdot \vec{\nabla} \vec{E}(\vec{r}_0) + \dots$$

$$- q \vec{E}(\vec{r}_0) + q \frac{\vec{a}}{2} \cdot \vec{\nabla} \vec{E}(\vec{r}_0) + \dots$$

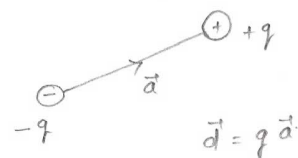
$$= \vec{a} \cdot \vec{\nabla} \vec{E}(\vec{r}_0)$$

$$= -\vec{a} \cdot \vec{\nabla} \vec{\nabla} \phi$$

$$= -\vec{\nabla} (\vec{a} \cdot \vec{\nabla} \phi)$$

$$= \vec{\nabla} \vec{a} \cdot \vec{E}$$

$$= -\vec{\nabla} U_{\text{dipole}}$$



$$\vec{r}_+ = \vec{r}_0 + \frac{\vec{a}}{2}$$

$$\vec{r}_- = \vec{r}_0 - \frac{\vec{a}}{2}$$

$$\|\vec{a} \cdot \vec{\nabla} \vec{E}\| \ll \|\vec{E}\|$$

$$a \ll \lambda$$

$\lambda \rightarrow$ characteristic wavelength.

(using $\vec{E} = -\vec{\nabla} \phi$ — electrostatic)

$$U_E = -\vec{d} \cdot \vec{E}$$

⑦ Torque on electric dipole

$$\begin{aligned}
 \vec{\tau}_E &= \sum_a (\vec{r}_a - \vec{r}_0) \times \vec{F}(\vec{r}_a) \\
 &= q \frac{\vec{a}}{2} \times \vec{E}(\vec{r}_1) - q \left(-\frac{\vec{a}}{2}\right) \times \vec{E}(\vec{r}_2) \\
 &= q \frac{\vec{a}}{2} \times \vec{E}(\vec{r}_0) + q \frac{\vec{a}}{2} \times \vec{E}(\vec{r}_0) \\
 &= \vec{d} \times \vec{E}
 \end{aligned}$$

⑧ Magnetic dipole moment

$$\vec{\mu} = \frac{1}{2} q \vec{r}_a \times \vec{v}_a = \frac{q}{2m} \vec{L}$$

$$U_B = - \vec{\mu} \cdot \vec{B}$$

$$\vec{\tau}_B = \vec{\mu} \times \vec{B}$$



$$\vec{\mu} = I A \hat{n}$$

$$\begin{aligned}
 \mu &= \frac{1}{2} q r v \\
 &= \frac{1}{2} \frac{q}{\Delta t} r \left(\frac{\Delta x}{2\pi r} \right) 2\pi r \\
 &= I \pi r^2 \\
 &= I A
 \end{aligned}$$