Homework No. 05 (Spring 2014)

PHYS 520B: Electromagnetic Theory

Due date: Friday, 2014 Mar 28, 4.30pm

1. From Maxwell's equations, without introducing potentials, show that the electric and magnetic fields satisfy the inhomogeneous wave equations

$$\left(-\nabla^2 + \frac{1}{c^2}\frac{\partial^2}{\partial t^2}\right)\mathbf{E}(\mathbf{r}, t) = -\frac{1}{\varepsilon_0}\nabla\rho(\mathbf{r}, t) - \frac{1}{\varepsilon_0}\frac{1}{c^2}\frac{\partial}{\partial t}\mathbf{J}(\mathbf{r}, t), \tag{1a}$$

$$\left(-\nabla^2 + \frac{1}{c^2}\frac{\partial^2}{\partial t^2}\right)\mathbf{B}(\mathbf{r}, t) = \mu_0 \mathbf{\nabla} \times \mathbf{J}(\mathbf{r}, t). \tag{1b}$$

Extend this result to magnetic charges and currents.

2. The 4-dimensional Euclidean Green's function satisfies

$$-\left(\frac{\partial^2}{\partial x_1^2} + \dots + \frac{\partial^2}{\partial x_4^2}\right) G_E(x_1, \dots, x_4) = \delta(x_1) \dots \delta(x_4). \tag{2}$$

(a) Show that the solution to this equation can be written as the Fourier transform

$$G_E(x_1, \dots, x_4) = \int_{-\infty}^{\infty} \frac{dk_1}{2\pi} \cdots \int_{-\infty}^{\infty} \frac{dk_4}{2\pi} \frac{e^{i(k_1x_1 + \dots + k_4x_4)}}{k_1^2 + \dots + k_4^2}.$$
 (3)

(b) Verify the integral

$$\frac{1}{M} = \int_0^\infty ds \, e^{-sM}.\tag{4}$$

(c) Using Eq. (4) in Eq. (3) show that

$$G_E(x_1, \dots, x_4) = \int_0^\infty ds \prod_{m=1}^4 \left[\int_{-\infty}^\infty \frac{dk_m}{2\pi} e^{-sk_m^2 + ik_m x_m} \right].$$
 (5)

(d) Show that

$$\int_{-\infty}^{\infty} dk_m e^{-sk_m^2 + ik_m x_m} = \sqrt{\frac{\pi}{s}} e^{-\frac{x_m^2}{4s}}$$
 (6)

(e) Using the integral of Eq. (6) in Eq. (5) and using the integral representation of Gamma function,

$$\Gamma(z) = \int_0^\infty \frac{ds}{s} s^z e^{-s},\tag{7}$$

show that

$$G_E(x_1, \dots, x_4) = \frac{1}{4\pi^2} \frac{1}{x_1^2 + \dots + x_4^2}.$$
 (8)

(f) By making the complex replacement

$$x_4 \to ict \equiv \lim_{\varepsilon \to 0+} e^{i(\frac{\pi}{2} - \varepsilon)} ct,$$
 (9)

show that

$$D_{+}(\mathbf{r},t) = iG_E(\mathbf{r},ict) \tag{10}$$

satisfies the differential equation

$$\left(-\nabla^2 + \frac{1}{c^2}\frac{\partial^2}{\partial t^2}\right)D_+(\mathbf{r}, t) = \delta^{(3)}(\mathbf{r})\delta(ct),\tag{11}$$

where $\mathbf{r} = (x_1, x_2, x_3)$, with the corresponding solution

$$D_{+}(\mathbf{r},t) = \lim_{\varepsilon \to 0+} \frac{i}{4\pi^2} \frac{1}{\mathbf{r}^2 - (ct)^2 + i\varepsilon'},\tag{12}$$

where $\varepsilon' = (ct)^2 \varepsilon$

(g) Using the δ -function representation

$$\pi\delta(x) = \lim_{\varepsilon \to 0+} \frac{\varepsilon}{x^2 + \varepsilon^2} = \lim_{\varepsilon \to 0+} \operatorname{Im} \frac{1}{x - i\varepsilon}$$
 (13)

and the identity

$$\delta(r^2 - (ct)^2) = \frac{1}{2r} \left[\delta(r - ct) + \delta(r + ct) \right]$$
(14)

show that

$$\operatorname{Re} D_{+}(\mathbf{r}, t) = \frac{1}{2} \left[\frac{\delta(r - ct)}{4\pi r} + \frac{\delta(r + ct)}{4\pi r} \right], \tag{15}$$

where the two terms here are the retarded and advanced Green's functions, respectively, up to numerical factors.

- (h) Refer problem 31.9 in Schwinger et al. for further discussion on this subject. (Will not be graded.)
- 3. Using the identity

$$\delta(F(x)) = \sum_{r} \frac{\delta(x - a_r)}{\left|\frac{dF}{dx}\right|_{x = a_r}},\tag{16}$$

where the sum on r runs over the roots a_r of the equation F(x) = 0, evaluate

$$\delta(\sin x)$$
, $\delta(\cos x)$, and $\delta(\tan x)$. (17)