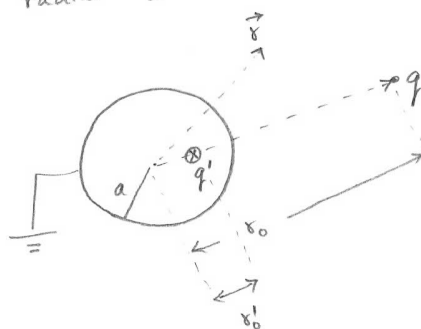


Method of images - Conducting sphere

- ① Consider a charge q outside a perfectly conducting sphere of radius a



- ② Electric field is zero inside a conductor, and can have only components normal to the surface. Any constant electric potential will achieve this. A grounded conductor means that this constant is chosen to be zero. Thus, we have.

$$\phi(\vec{r}) = 0 \quad (\text{grounded conductor.})$$

- ③ Using symmetry arguments we conclude:
 (i) Image charge is on the radial line.
 (ii) Image charge is opposite in sign.

④
$$\phi(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{q}{|\vec{r} - \vec{r}_0|} - \frac{1}{4\pi\epsilon_0} \frac{q'}{|\vec{r} - \vec{r}'|}$$

$$\textcircled{5} \quad \phi(\vec{a}) = 0$$

$$\vec{r} \cdot \vec{r}_0 = 2ar_0 \cos r$$

$$\Rightarrow \frac{q}{|\vec{a} - \vec{r}_0|} - \frac{q'}{|\vec{a} - \vec{r}'_0|} = 0$$

$$\frac{q}{\sqrt{a^2 + r_0^2 - 2ar_0 \cos r}} = \frac{q'}{\sqrt{a^2 + r_0'^2 - 2ar_0' \cos r}}$$

$$q^2 (a^2 + r_0'^2 - 2ar_0' \cos r) = q'^2 (a^2 + r_0^2 - 2ar_0 \cos r)$$

$$q^2 (a^2 + r_0'^2) - q'^2 (a^2 + r_0^2) + 2a(r_0 q'^2 - r_0' q^2) \cos r = 0$$

⑥ Since the above equation is true for all r , we have.

$$q^2 (a^2 + r_0'^2) = q'^2 (a^2 + r_0^2) \quad \text{--- (i)}$$

$$\text{and} \quad r_0 q'^2 - r_0' q^2 = 0 \quad \text{--- (ii)}$$

⑦ Thus, we have.

$$\frac{q'^2}{q^2} = \frac{r_0'}{r_0}$$

⑧ Using ⑦ in ⑥ - (i)

$$r_0 (a^2 + r_0'^2) = r_0' (a^2 + r_0^2)$$

$$\left(\frac{r_0'}{r_0}\right)^2 - \frac{r_0'}{r_0} \left(1 + \frac{a^2}{r_0^2}\right) + \frac{a^2}{r_0^2} = 0$$

$$\frac{r_0'}{r_0} = \frac{1}{2} \left(1 + \frac{a^2}{r_0^2}\right) \pm \frac{1}{2} \sqrt{\left(1 + \frac{a^2}{r_0^2}\right)^2 - 4 \frac{a^2}{r_0^2}}$$

$$= \frac{1}{2} \left(1 + \frac{a^2}{r_0^2}\right) \pm \frac{1}{2} \left(1 - \frac{a^2}{r_0^2}\right) = \begin{cases} 1 \\ \frac{a^2}{r_0^2} \end{cases} \rightarrow \text{trivial.}$$

⑨ Using ⑧ in ⑦ we have.

$$\frac{q'}{q^2} = \frac{r_0'}{r_0} = \frac{a^2}{r_0^2}$$

⑩ Using ⑨ in ④

$$\begin{aligned} \phi(\vec{r}) &= \frac{1}{4\pi\epsilon_0} \frac{q}{\sqrt{r^2 + r_0^2 - 2rr_0\cos r}} - \frac{1}{4\pi\epsilon_0} \frac{q'}{\sqrt{r^2 + r_0'^2 - 2rr_0'\cos r}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{\sqrt{r^2 + r_0^2 - 2rr_0\cos r}} - \frac{1}{4\pi\epsilon_0} \frac{\frac{a^2}{r_0} q}{\sqrt{r^2 + \frac{a^4}{r_0^2} - 2rr_0\frac{a^2}{r_0^2}\cos r}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{\sqrt{r^2 + r_0^2 - 2rr_0\cos r}} - \frac{1}{4\pi\epsilon_0} \frac{q}{\sqrt{r^2 \frac{r_0^2}{a^2} + a^2 - 2rr_0\cos r}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{\sqrt{r^2 + r_0^2 - 2rr_0\cos r}} - \frac{1}{4\pi\epsilon_0} \frac{q}{\sqrt{\left(\frac{r}{a}r_0\right)^2 + a^2 - 2rr_0\cos r}} \end{aligned}$$

which is clearly zero at $r = a$.

- ⑪ Charge density on the surface of a perfect conductor:

$$\vec{\nabla} \cdot (\epsilon_0 \vec{E}) = \rho - \vec{\nabla} \cdot \vec{P}$$

⑫ Let $\vec{P} = \vec{P}_0 \theta(R-r)$

$$\begin{aligned} \vec{\nabla} \cdot \vec{P} &= (\vec{\nabla} \cdot \vec{P}_0) \theta(R-r) + \vec{P}_0 \cdot \vec{\nabla} \theta(R-r) \\ &= (\vec{\nabla} \cdot \vec{P}_0) \theta(R-r) + \underbrace{(\vec{P}_0 \cdot \hat{r}) \delta(r-R)}_{\nabla(\theta, \phi)} \end{aligned}$$

↓
this term is necessarily zero for a perfect conductor

- ⑬ In general, we will have

$$\rho_{\text{eff}}(\vec{r}) = -\vec{\nabla} \cdot \vec{P} = +\nabla(\theta, \phi) \delta(\xi - \xi_0)$$

- ξ is the coordinate of the surface of conductor.
→ For a conductor, we can only consider one side of surface, unless it is a infinitesimally thin conductor.

- Examples: Surface charge density on a conducting plane $\left. \vphantom{\begin{array}{l} \text{Surface charge density} \\ \text{on a conducting plane} \end{array}} \right\} = +\nabla(x, y) \delta(z - z_0)$
Surface charge density on a conducting sphere $\left. \vphantom{\begin{array}{l} \text{Surface charge density} \\ \text{on a conducting sphere} \end{array}} \right\} = +\nabla(\theta, \phi) \delta(r - r_0)$

(14) For a conducting sphere we have.

$\perp \rightarrow$ perpendicular to normal.

$$\vec{\nabla} \cdot \epsilon_0 \vec{E} = + \nabla(\theta, \phi) \delta(r-a)$$

$$\epsilon_0 \frac{\partial}{\partial r} E_r + \epsilon_0 \underbrace{\vec{\nabla}_\perp \cdot \vec{E}_\perp}_{=0 \text{ on the conductor}} = + \nabla(\theta, \phi) \delta(r-a)$$

$$\epsilon_0 \frac{\partial}{\partial r} E_r = + \nabla(\theta, \phi) \delta(r-a)$$

Integrating around $r=a$, we have.

$$\epsilon_0 \int_{r=a-\delta}^{r=a+\delta} dr \frac{\partial}{\partial r} E_r = + \nabla(\theta, \phi)$$

$$E_r(a+\delta) - E_r(a-\delta) = + \frac{1}{\epsilon_0} \nabla(\theta, \phi)$$

$\downarrow$
 $=0$ inside a conductor

$$\begin{aligned} \nabla(\theta, \phi) &= + \epsilon_0 E_r(a+\delta) \\ &= - \epsilon_0 \frac{\partial}{\partial r} \phi \Big|_{r=a+\delta} \end{aligned}$$

(15) Using (10) in (14)

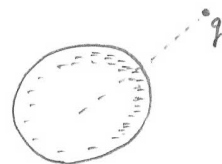
$$\begin{aligned} \nabla(\theta, \phi) &= - \frac{1}{4\pi} \frac{\partial}{\partial r} \left[\frac{q}{\sqrt{r^2 + r_0^2 - 2rr_0 \cos \Gamma}} - \frac{q}{\sqrt{r^2 \frac{r_0^2}{a^2} + a^2 - 2rr_0 \cos \Gamma}} \right] \Big|_{r=a+b} \\ &= - \frac{q}{4\pi} \left[\frac{(r_0 \cos \Gamma - r)}{(r^2 + r_0^2 - 2rr_0 \cos \Gamma)^{3/2}} - \frac{(r_0 \cos \Gamma - r \frac{r_0^2}{a^2})}{(r^2 \frac{r_0^2}{a^2} + a^2 - 2rr_0 \cos \Gamma)^{3/2}} \right] \Big|_{r=a+b} \\ &= - \frac{q}{4\pi} \frac{a \left(\frac{r_0^2}{a^2} - 1 \right)}{(a^2 + r_0^2 - 2ar_0 \cos \Gamma)^{3/2}} \end{aligned}$$

$$\begin{aligned}
 (16) \quad \nabla(\theta, \phi) &= -\frac{q}{4\pi a^2} \frac{\left(\frac{r_0^2}{a^2} - 1\right)}{\left(1 + \frac{r_0^2}{a^2} - 2\frac{r_0}{a} \cos \theta\right)^{\frac{3}{2}}} \\
 &= -\frac{q}{4\pi a^2} \frac{\frac{a}{r_0} \left(1 - \frac{a^2}{r_0^2}\right)}{\left(1 + \frac{a^2}{r_0^2} - 2\frac{a}{r_0} \cos \theta\right)^{\frac{3}{2}}}
 \end{aligned}$$

$$\begin{aligned}
 (17) \quad \nabla(0, \phi) &= -\frac{q}{4\pi a^2} \frac{\frac{a}{r_0} \left(1 - \frac{a^2}{r_0^2}\right)}{\left(1 - \frac{a}{r_0}\right)^3} = -\frac{q}{4\pi a^2} \frac{\frac{a}{r_0} \left(1 + \frac{a}{r_0}\right)}{\left(1 - \frac{a}{r_0}\right)^2} \\
 \nabla(\pi, \phi) &= -\frac{q}{4\pi a^2} \frac{\frac{a}{r_0} \left(1 - \frac{a^2}{r_0^2}\right)}{\left(1 + \frac{a}{r_0}\right)^3} = -\frac{q}{4\pi a^2} \frac{\frac{a}{r_0} \left(1 - \frac{a}{r_0}\right)}{\left(1 + \frac{a}{r_0}\right)^2}
 \end{aligned}$$

(18) Note that,

$$\frac{1 + \frac{a}{r_0}}{\left(1 - \frac{a}{r_0}\right)^2} > 1 \quad \text{and} \quad \frac{1 - \frac{a}{r_0}}{\left(1 + \frac{a}{r_0}\right)^2} < 1$$



(19) To do: Multipole expansion of $\nabla(\theta, \phi)$.

charge q along z -axis.

$$\begin{aligned}
 (20) \quad \nabla_{\text{tot}} &= a^2 \int_0^{2\pi} d\phi \int_0^\pi \sin \theta d\theta \nabla(\theta, \phi) \\
 &= 2\pi a^2 \int_0^\pi \sin \theta d\theta \nabla(\theta, \phi) \\
 &= 2\pi a^2 (-1) \frac{q}{4\pi a^2} \frac{\frac{a}{r_0} \left(1 - \frac{a^2}{r_0^2}\right)}{\left(1 + \frac{a^2}{r_0^2}\right)^{3/2}} \int_0^\pi \sin \theta d\theta \frac{1}{\left[1 - \frac{2\frac{a}{r_0}}{\left(1 + \frac{a^2}{r_0^2}\right)} \cos \theta\right]^{\frac{3}{2}}}
 \end{aligned}$$

$$(21) \quad \gamma_{tot} = - \frac{q}{2} \frac{\frac{a}{\gamma_0} \left(1 - \frac{a^2}{\gamma_0^2}\right)}{\left(1 + \frac{a^2}{\gamma_0^2}\right)^{3/2}} \int_{-1}^1 dt \frac{1}{\left[1 - \left(\frac{2 \frac{a}{\gamma_0}}{1 + \frac{a^2}{\gamma_0^2}}\right) t\right]^{3/2}}$$

$$(22) \quad \int_{-1}^1 dt \frac{1}{(1-xt)^{3/2}} = - \int_{1-x}^{1+x} \frac{ds}{s} \frac{1}{s^{3/2}} \quad \begin{array}{l} 1-xt = s \\ -x dt = ds \end{array}$$

$$= \frac{1}{x} \int_{1-x}^{1+x} \frac{ds}{s^{3/2}}$$

$$= \frac{1}{x} (-2) \frac{1}{\sqrt{s}} \Big|_{s=1-x}^{s=1+x} = - \frac{2}{x} \left[\frac{1}{\sqrt{1+x}} - \frac{1}{\sqrt{1-x}} \right]$$

$$(23) \quad \gamma_{tot} = - \frac{q}{2} \frac{\frac{a}{\gamma_0} \left(1 - \frac{a^2}{\gamma_0^2}\right)}{\left(1 + \frac{a^2}{\gamma_0^2}\right)^{3/2}} (-1) \frac{\cancel{2} \left(1 + \frac{a^2}{\gamma_0^2}\right)}{\cancel{2} \frac{a}{\gamma_0}} \left[\frac{1}{\sqrt{1 + \left(\frac{2 \frac{a}{\gamma_0}}{1 + \frac{a^2}{\gamma_0^2}}\right)}} - \frac{1}{\sqrt{1 - \left(\frac{2 \frac{a}{\gamma_0}}{1 + \frac{a^2}{\gamma_0^2}}\right)}} \right]$$

$$= \frac{q}{2} \frac{\left(1 - \frac{a^2}{\gamma_0^2}\right)}{\left(1 + \frac{a^2}{\gamma_0^2}\right)^{1/2}} \left[\frac{\sqrt{1 + \frac{a^2}{\gamma_0^2}}}{\left(1 + \frac{a}{\gamma_0}\right)} - \frac{\sqrt{1 + \frac{a^2}{\gamma_0^2}}}{\left(1 - \frac{a}{\gamma_0}\right)} \right]$$

$$= \frac{q}{2} \left(1 - \frac{a^2}{\gamma_0^2}\right) \left[\frac{1}{\left(1 + \frac{a}{\gamma_0}\right)} - \frac{1}{\left(1 - \frac{a}{\gamma_0}\right)} \right]$$

$$= \frac{q}{2} \left(1 - \frac{a^2}{\gamma_0^2}\right) \frac{(-2 \frac{a}{\gamma_0})}{\left(1 - \frac{a^2}{\gamma_0^2}\right)}$$

$$= - q \frac{a}{\gamma_0}$$

(24)

$$q' = -q \frac{a}{r_0}$$

$$r_0' = r_0 \frac{a^2}{r_0^2}$$



$$F = \frac{1}{4\pi\epsilon_0} \frac{q q'}{|r_0 - r_0'|^2} = - \frac{q^2}{4\pi\epsilon_0} \frac{\frac{a}{r_0}}{\left|r_0 - r_0 \frac{a^2}{r_0^2}\right|^2} = - \frac{q^2}{4\pi\epsilon_0} \frac{a r_0}{(r_0^2 - a^2)^2}$$

$$E_{int} = \frac{1}{2} \frac{1}{4\pi\epsilon_0} \frac{q q'}{|r_0 - r_0'|} = - \frac{q^2}{4\pi\epsilon_0} \frac{1}{2} \frac{\frac{a}{r_0}}{\left|r_0 - r_0 \frac{a^2}{r_0^2}\right|} = - \frac{1}{2} \frac{q^2}{4\pi\epsilon_0} \frac{a}{(r_0^2 - a^2)}$$

$$F = - \frac{\partial}{\partial r_0} E_{int}$$

$$= - \frac{\partial}{\partial r_0} \left(\frac{-1}{2} \frac{q^2}{4\pi\epsilon_0} \frac{a}{(r_0^2 - a^2)} \right)$$

$$= \frac{1}{2} \frac{q^2}{4\pi\epsilon_0} (-1) \frac{a}{(r_0^2 - a^2)^2} \cdot 2 r_0$$

$$= - \frac{q^2}{4\pi\epsilon_0} \frac{a r_0}{(r_0^2 - a^2)^2} \quad \checkmark$$