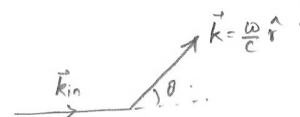


Light scattering off polymers

- ① Scattering amplitude for light scattering off a polarizable atom/molecule in the Rayleigh approximation is.

$$\vec{E}_0(\theta, \phi) = -\alpha_0 \frac{\omega^2}{c^2} \hat{r} \times (\hat{r} \times \vec{E}_0) e^{-i(\vec{k} - \vec{k}_{in}) \cdot \vec{r}_0}$$



$\vec{r}_0$ is the position of the molecule.

- ② We found that

$$|\vec{E}_0(\theta, \phi)|^2 = \alpha^2 \frac{\omega^4}{c^4} E_0^2 \begin{cases} 1, & \perp \text{ pol.}, \\ \cos^2 \theta, & \parallel \text{ pol.}, \\ \frac{1 + \cos^2 \theta}{2}, & \text{unpolarized light.} \end{cases}$$

- ③ For scattering off N particles we have

$$|F(\theta, \phi)|^2 = \left| \sum_{m=1}^N F_m(\theta, \phi) \right|^2,$$

which leads to interference effects. Crystallography involves the case when the N molecules are orderly positioned. Here we shall study scattering off a polymer chain whose position is known as a probability distribution.

④ Polymer is, to the first approximation, modelled as a Brownian motion in 3+1 space-time dimensions.
Let us consider Brownian motion in 1+1 space-time dimensions:

a - length of each step → starts at origin.

x - distance from origin.

s - total distance travelled.

$P(x, s)$ - probability of finding the particle at x at time t . Note s is a measure of t .

⑤ We observe that

$$P(x, s+a) = \frac{1}{2} [P(x+a, s) + P(x-a, s)]$$

Taylor expanding and keeping only the leading terms,

$$P(x, s) + a \frac{\partial}{\partial s} P(x, s) = \frac{1}{2} \left[P(x, s) + a \frac{\partial}{\partial x} P(x, s) + \frac{a^2}{2} \frac{\partial^2}{\partial x^2} P(x, s) + \dots \right. \\ \left. + P(x, s) - a \frac{\partial}{\partial x} P(x, s) + \frac{a^2}{2} \frac{\partial^2}{\partial x^2} P(x, s) + \dots \right]$$

$$a \frac{\partial}{\partial s} P(x, s) = \frac{a^2}{2} \frac{\partial^2}{\partial x^2} P(x, s)$$

This is the diffusion equation in 1+1 dimension.

- ⑥ This idea can be generalized to 3+1 dimension
to derive the diffusion equation in 3+1 dimen.

$$\frac{\partial}{\partial s} P(\vec{r}, s) = \frac{a^2}{6} \nabla^2 P(\vec{r}, s)$$

For the initial condition

$$P(\vec{r}, 0) = \delta^{(3)}(\vec{r})$$

we have the solution

$$P(\vec{r}, s) = \left(\frac{3}{2\pi a s} \right)^{\frac{3}{2}} e^{-\frac{3r^2}{2\pi a s}}$$

which satisfies

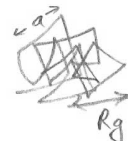
$$\int d^3r P(\vec{r}, s) = 1.$$

- ⑦ The quantity

$$\langle r^2 \rangle = \int d^3r r^2 P(\vec{r}, s)$$

$$= a s$$

is a measure of the square of polymer radius.



- ⑧ Let

$$L = Na.$$

$$R^2 = \langle r^2 \rangle = Na^2 \Rightarrow R = \sqrt{Na}.$$

$$R_g^2 = \frac{1}{6} Na^2 \rightarrow \text{gyration radius.}$$

⑨ The scattering amplitude off each monomer in the polymer is

$$\vec{q} = \vec{k} - \vec{k}_{in}$$

$$F_m(\theta, \phi) = -\alpha(\vec{R}_m) \frac{\omega^2}{c^2} \hat{r} \times (\hat{r} \times \vec{E}_0) e^{-i \vec{q} \cdot \vec{R}_m}$$

$\vec{R}_m$ - position of monomer

Let $\alpha(\vec{R}_m) = \alpha$

⑩ Total scattering amplitude is

$$F(\theta, \phi) = \sum_{m=1}^N F_m(\theta, \phi)$$

$$= -\alpha \frac{\omega^2}{c^2} \hat{r} \times (\hat{r} \times \vec{E}_0) \sum_{m=1}^N e^{-i \vec{q} \cdot \vec{R}_m}$$

$$|F(\theta, \phi)|^2 = \alpha^2 \frac{\omega^4}{c^4} |\hat{r} \times (\hat{r} \times \vec{E}_0)|^2 \sum_{m=1}^N \sum_{n=1}^N e^{i \vec{q} \cdot (\vec{R}_n - \vec{R}_m)}$$

⑪

Scattering cross section

$$= \alpha^2 \frac{\omega^4}{c^4} E_0^2 \frac{1 + \cos^2 \theta}{2} \sum_{m=1}^N \sum_{n=1}^N e^{i \vec{q} \cdot (\vec{R}_n - \vec{R}_m)}$$

(12)

For a polymer the scattering cross section needs to be averaged with the probability amplitude.

$$\begin{aligned}
 \langle |F(\theta, \phi)|^2 \rangle &= \alpha^2 \frac{\omega^4}{c^4} E_0^2 \frac{1 + \cos^2 \theta}{2} \\
 &\times \sum_{m=1}^N \sum_{n=1}^N \int d^3(\vec{R}_m - \vec{R}_n) P(\vec{R}_m - \vec{R}_n) e^{i\vec{q} \cdot (\vec{R}_n - \vec{R}_m)} \\
 &= \alpha^2 \frac{\omega^4}{c^4} E_0^2 \frac{1 + \cos^2 \theta}{2} \sum_{m=1}^N \sum_{n=1}^N \int d^3(\vec{R}_m - \vec{R}_n) \left(\frac{3}{2\pi |n-m| a^2} \right)^{\frac{3}{2}} e^{-\frac{3(\vec{R}_n - \vec{R}_m)^2}{2|n-m| a^2}} e^{i\vec{q} \cdot (\vec{R}_n - \vec{R}_m)} \\
 &= \alpha^2 \frac{\omega^4}{c^4} E_0^2 \frac{1 + \cos^2 \theta}{2} \sum_{m=1}^N \sum_{n=1}^N \left(\frac{1}{\sqrt{\pi}} \right)^3 \int d^3 \vec{t} e^{-t^2 + i\vec{a} \cdot \vec{t}} e^{i\vec{q} \cdot \sqrt{\frac{2|n-m| a^2}{3}} \vec{t}} \\
 &= \alpha^2 \frac{\omega^4}{c^4} E_0^2 \frac{1 + \cos^2 \theta}{2} \sum_{m=1}^N \sum_{n=1}^N \left(\frac{1}{\sqrt{\pi}} \right)^3 \int d^3 \vec{t} e^{-(\vec{t} - i\frac{\vec{a}}{2})^2} e^{(i\frac{\vec{a}}{2})^2} \\
 &= \alpha^2 \frac{\omega^4}{c^4} E_0^2 \frac{1 + \cos^2 \theta}{2} \sum_{m=1}^N \sum_{n=1}^N e^{-\frac{b^2}{4}} \\
 &= \alpha^2 \frac{\omega^4}{c^4} E_0^2 \frac{1 + \cos^2 \theta}{2} \sum_{m=1}^N \sum_{n=1}^N e^{-\frac{1}{6} q^2 a^2 |n-m|}
 \end{aligned}$$

(13)
$$\sum_{m=1}^N \sum_{n=1}^N e^{-\frac{1}{6} q^2 a^2 |n-m|} = \int_0^L \frac{dx}{a} \int_0^L \frac{dx'}{a} e^{-\frac{1}{6} q^2 a |x-x'|}$$

$x = na$
 $x' = ma$

$t = \frac{1}{6} q^2 a x$
 $dt = \frac{q^2 a^2}{6} \frac{dx}{a}$

$$= \left(\frac{6}{q^2 a^2}\right)^2 \int_0^{\frac{1}{6} q^2 a L} dt \int_0^{\frac{1}{6} q^2 a L} dt' e^{-|t-t'|}$$

(14)
$$\begin{aligned} \int_0^c dt \int_0^c dt' e^{-|t-t'|} &= \int_0^c dt \int_0^t dt' e^{-(t-t')} + \int_0^c dt \int_t^c dt' e^{-(t'-t)} \\ &= \int_0^c dt e^{-t} (e^t - 1) + \int_0^c dt e^t (\bar{e}^{-t} - \bar{e}^{-c}) \\ &= \int_0^c dt (1 - \bar{e}^t) + \int_0^c dt (1 - \bar{e}^c e^t) \\ &= \frac{c}{1} + (\bar{e}^{-c} - 1) + \frac{c}{1} - \bar{e}^c (e^c - 1) \\ &= 2c - 2 + 2\bar{e}^c \\ &= 2(\bar{e}^c + c - 1) \end{aligned}$$

$R_g^2 = \frac{aL}{6}$
 $= \frac{1}{6} Na^2$

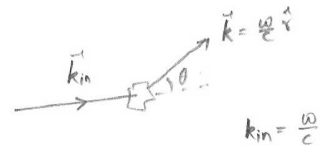
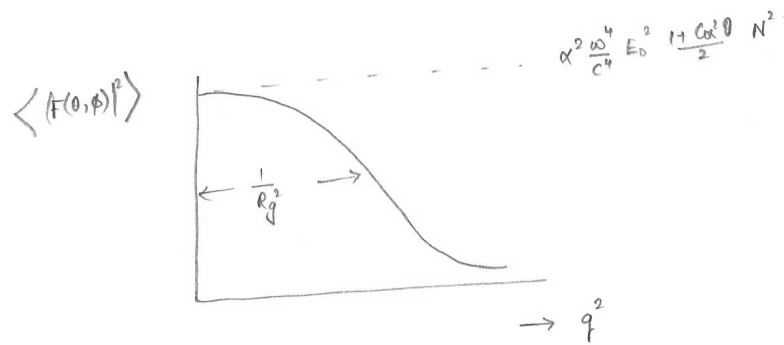
(15) Using (14) in (13)

$$\begin{aligned} \sum_{m=1}^N \sum_{n=1}^N e^{-\frac{1}{6} q^2 a^2 |n-m|} &= \left(\frac{6}{q^2 a^2}\right)^2 2 \left[e^{-\frac{q^2 a L}{6}} + \frac{q^2 a L}{6} - 1 \right] \\ &= \left(\frac{N}{q^2 R_g^2}\right)^2 2 \left[e^{-q^2 R_g^2} + q^2 R_g^2 - 1 \right] \end{aligned}$$

(16) Using (15) in (12) we have.

$$\langle |F(\theta, \phi)|^2 \rangle = \alpha^2 \frac{\omega^4}{c^4} E_0^2 \frac{1 + \cos^2 \theta}{2} N^2 \frac{2}{(q^2 R_g^2)^2} \left[e^{-q^2 R_g^2} + q^2 R_g^2 - 1 \right]$$

(17)
$$\frac{2}{x^4} (e^{-x^2} + x^2 - 1) = \begin{cases} 1 & x \ll 1. \\ \frac{2}{x^2} & 1 \ll x. \end{cases}$$



Note that

$$\begin{aligned} q^2 &= (\vec{k} - \vec{k}_{in})^2 \\ &= k^2 + k_{in}^2 - 2 \vec{k} \cdot \vec{k}_{in} \\ &= 2k^2 (1 - \cos \theta) \\ &= 4k^2 \sin^2 \frac{\theta}{2} \\ &= 4 \frac{\omega^2}{c^2} \sin^2 \frac{\theta}{2} \end{aligned}$$

Thus, angular variation is equivalent to varying q^2 .