

- ① Maxwell's equations , for electrostatics , reduces to finding the electric potential ,
- $$-\vec{\nabla} \cdot [\epsilon(\vec{r}) \vec{\nabla} \phi(\vec{r})] = \rho(\vec{r}) .$$
- ↓
dielectric material .
- ↳ charge distribution

- ② The Green function corresponding to this is
- $$-\vec{\nabla} \cdot [\epsilon(\vec{r}) \vec{\nabla} G(\vec{r}, \vec{r}')] = \delta^{(3)}(\vec{r} - \vec{r}')$$

- ③ For planar geometry ,
- $$\epsilon(\vec{r}) = \epsilon(z),$$

we have.

$$\phi(\vec{r}) = \int d^3r' G(\vec{r}, \vec{r}') \rho(\vec{r}')$$

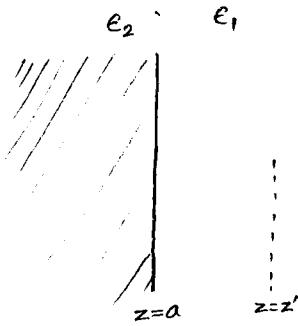
$$G(\vec{r}, \vec{r}') = \int \frac{d^2k_\perp}{(2\pi)^2} e^{i\vec{k}_\perp \cdot (\vec{r} - \vec{r}')_\perp} g_e(z, z'; k_\perp)$$

Greens function $g_e(z, z'; k_\perp)$ satisfies

where the reduced

$$\left[-\frac{\partial}{\partial z} \epsilon(z) \frac{\partial}{\partial z} + k_\perp^2 \epsilon(z) \right] g_e(z, z'; k_\perp) = \delta(z - z') .$$

④ Let us consider the situation



$$e(z) = \begin{cases} e_2, & z < a, \\ e_1, & z > a. \end{cases}$$

⑤ For say, $a < z'$, we can look for solutions of the kind, because in each of the following regions the $e(z)$ is uniform and the δ -function does not contribute. Thus,

$$g_e(z, z'; k_1) = \begin{cases} A e^{k_1 z} + B e^{-k_1 z}, & z < a < z', \\ C e^{k_1 z} + D e^{-k_1 z}, & a < z < z', \\ E e^{k_1 z} + F e^{-k_1 z}, & a < z' < z. \end{cases}$$

⑥ The constraints that there solutions have to satisfy are -

$$(i) \quad g(-\infty, z'; k_1) = 0$$

$$(ii) \quad g(+\infty, z'; k_1) = 0$$

$$(iii) \quad g(z, z'; k_1) \Big|_{z=z'-\delta}^{z=z'+\delta} = 0$$

$$(iv) \quad e(z) \frac{\partial}{\partial z} g(z, z'; k_1) \Big|_{z=z'-\delta}^{z=z'+\delta} = -1 \quad \rightarrow \text{notice the } e(z) \text{ factor.}$$

$$(v) \quad g(z, z'; k_1) \Big|_{z=a-\delta}^{z=a+\delta} = 0$$

$$(vi) \quad e(z) \frac{\partial}{\partial z} g(z, z'; k_1) \Big|_{z=a-\delta}^{z=a+\delta} = 0$$

⑦ The six conditions can be used to determine the six unknowns in ⑤, as will be described in

homework, to obtain

$$g_e(z, z'; k_1) = \begin{cases} \frac{1}{e_2} \frac{1}{2k_1} e^{-k_1|z-z'|} + \frac{1}{e_2} \frac{1}{2k_1} \frac{e_2 - e_1}{e_2 + e_1} e^{-k_1|z-a|} e^{-k_1|z'-a|}, & z' < a \\ \frac{1}{e_1} \frac{1}{2k_1} e^{-k_1|z-z'|} + \frac{1}{e_1} \frac{1}{2k_1} \frac{e_1 - e_2}{e_1 + e_2} e^{-k_1|z-a|} e^{-k_1|z'-a|}, & a < z' \end{cases}$$

⑧ Interpreting the second term as the contribution due to a "image charge" we identify the magnitude of the image charge as.

$$q_{\text{image}} = \begin{cases} \frac{\epsilon_2 - \epsilon_1}{\epsilon_2 + \epsilon_1}, & \text{if } z' < a, \\ \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2}, & \text{if } a < z'. \end{cases}$$

and the position of the image to be

$$|\vec{r} - \vec{r}'_{\text{image}}| = (x - x') \hat{i} + (y - y') \hat{j} + |z - a| + |z' - a| \hat{k}$$

or

$$|z - z'_{\text{image}}| = |z - a| + |z' - a|,$$

which implies.

$$z'_{\text{image}} = \begin{cases} 2a - z', & \text{if } z' < a, z < a, \\ z', & \text{if } z' < a, a < z, \\ z', & \text{if } a < z', z < a, \\ 2a - z', & \text{if } a < z', a < z'. \end{cases}$$

Thus, the image charge is located behind the "mirror"

if the observation point is on the same side as the charge.