

① For two charges we have the Coulomb interaction energy

$$U = \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{|\vec{r}_1 - \vec{r}_2|} \quad \text{--- } q_1, \vec{r}_1$$

where $\vec{r}_1$ and $\vec{r}_2$ are positions of the charges.

② The interaction energy for three charges is

$$U_{int} = \frac{1}{4\pi\epsilon_0} \left[\frac{q_1 q_2}{|\vec{r}_1 - \vec{r}_2|} + \frac{q_2 q_3}{|\vec{r}_2 - \vec{r}_3|} + \frac{q_1 q_3}{|\vec{r}_1 - \vec{r}_3|} \right]$$

$$= \frac{1}{2} \sum_{i=1}^3 \sum_{j=1, j \neq i}^3 \frac{1}{4\pi\epsilon_0} \frac{q_i q_j}{|\vec{r}_i - \vec{r}_j|}$$

where $i \neq j$: removes the self interaction terms

$\frac{1}{2}$: removes the double counting

③ For a continuous charge distribution the above generalizes to

$$U_{int} = \frac{1}{2} \int d^3\vec{r} \int d^3\vec{r}' \frac{1}{4\pi\epsilon_0} \frac{\rho(\vec{r}) \rho(\vec{r}')}{|\vec{r} - \vec{r}'|} \quad \text{--- } E_{self}$$

$$= \frac{1}{2} \int d^3\vec{r} \int d^3\vec{r}' \rho(\vec{r}) G_0(\vec{r}, \vec{r}') \rho(\vec{r}') \quad \text{--- } E_{self},$$

where E_{self} is the self energy, which does not contribute to force. Note that $G_0(\vec{r}, \vec{r}')$ is the free Green's function.

④ For example, for two point charges

$$\rho(\vec{r}) = q_1 \delta^{(3)}(\vec{r} - \vec{r}_1) + q_2 \delta^{(3)}(\vec{r} - \vec{r}_2)$$

we have

$$U_{int} = \frac{1}{2} \int d^3r \int d^3r' \frac{1}{4\pi\epsilon_0} [q_1 \delta^{(3)}(\vec{r} - \vec{r}_1) + q_2 \delta^{(3)}(\vec{r} - \vec{r}_2)] \frac{1}{|\vec{r} - \vec{r}'|} [q_1 \delta^{(3)}(\vec{r}' - \vec{r}_1) + q_2 \delta^{(3)}(\vec{r}' - \vec{r}_2)]$$

— Eself

$$= \frac{1}{4\pi\epsilon_0} \frac{q_1 q_2}{|\vec{r}_1 - \vec{r}_2|}$$

⑤ The interaction energy between a charge and a dielectric material is given by

$$U_{int} = \frac{1}{2} \int d^3r \int d^3r' \rho(\vec{r}) \left[G(\vec{r}, \vec{r}') - G_0(\vec{r}, \vec{r}') \right] \rho(\vec{r}')$$

$\downarrow$ Green's function including the dielectric material. $\rightarrow$ free Green's function

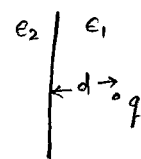
⑥ From earlier lectures we have the free Green's function

$$G_0(\vec{r}, \vec{r}') = \frac{1}{\epsilon_0} \int \frac{d^3k}{(2\pi)^3} e^{i\vec{k} \cdot (\vec{r} - \vec{r}')} \frac{1}{2k_z} e^{-k_z |z - z'|}$$

and the Green's function for a dielectric slab (occupying half space)

$$G(\vec{r}, \vec{r}') = \frac{1}{\epsilon_1} \int \frac{d^3k}{(2\pi)^3} e^{i\vec{k} \cdot (\vec{r} - \vec{r}')} \left[\frac{1}{2k_z} e^{-k_z |z - z'|} + \frac{1}{2k_z} \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2} e^{-k_z |z|} e^{-k_z |z'|} \right]$$

⑦ Thus, the interaction energy of a point charge and a dielectric plate is



$$\begin{aligned}
 U_{int} &= \frac{1}{2} \int d^3r \int d^3r' \rho(\vec{r}) [G(\vec{r}, \vec{r}') - G_0(\vec{r}, \vec{r}')] \rho(\vec{r}') \\
 &= \frac{1}{2} \int d^3r \int d^3r' \rho \delta^{(3)}(\vec{r} - \vec{r}_0) [G(\vec{r}, \vec{r}') - G_0(\vec{r}, \vec{r}')] \rho \delta^{(3)}(\vec{r}' - \vec{r}_0) \\
 &= \frac{q^2}{2} [G(\vec{r}_0, \vec{r}_0) - G_0(\vec{r}_0, \vec{r}_0)] \\
 &= \frac{q^2}{2} \frac{1}{\epsilon_1} \int \frac{d^3k}{(2\pi)^3} e^{i\vec{k}_1 \cdot 0} \frac{1}{2k_1} \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2} e^{-k_1 d} e^{-k_1 d} \\
 &= \frac{q^2}{16\pi^2} \frac{1}{\epsilon_1} \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2} \int_0^\infty k_1 dk_1 \underbrace{\int_0^{2\pi} d\phi}_{2\pi} \frac{1}{k_1} e^{-2k_1 d} \\
 &= \frac{q^2}{8\pi} \frac{1}{\epsilon_1} \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2} \frac{1}{2d} \underbrace{\int_0^\infty dx e^{-x}}_{=1} \\
 &= \frac{1}{4\pi\epsilon_1} \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2} \frac{q^2}{4d}
 \end{aligned}$$

⑧ $F = - \frac{\partial}{\partial d} U_{int} = \frac{1}{4\pi\epsilon_1} \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2} \frac{q^2}{(2d)^2}$

which is constant with the image interpretation:

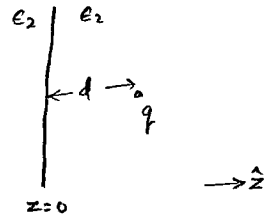
$$q_{image} = q \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2}$$

$$distance = 2d.$$

④ The force between the charge and the dielectric can also be calculated from the force on the induced electric charge on the dielectric plate,

$$F = \int_{-\infty}^{+\infty} dx \int_{-\infty}^{+\infty} dy \quad \sigma(x,y) E_z(x,y,0)$$

where $\sigma(x,y)$ is the charge induced in the dielectric due to the presence of q .



⑩ We identify the expression for the charge density from

$$\begin{aligned} \rho(\vec{r}) - \vec{\nabla} \cdot \vec{P} &= \rho - \vec{\nabla} \cdot [(\epsilon - \epsilon_0) \vec{E}] \\ &= \rho - \vec{\nabla} \cdot \left[\epsilon \vec{E} \left(1 - \frac{\epsilon_0}{\epsilon}\right) \right] \\ &= \rho - \rho \left(1 - \frac{\epsilon_0}{\epsilon}\right) - \epsilon \vec{E} \cdot \vec{\nabla} \left(1 - \frac{\epsilon_0}{\epsilon}\right) \\ &= \rho(\vec{r}) \frac{\epsilon_0}{\epsilon(\vec{r})} + \epsilon(\vec{r}) \vec{E} \cdot \vec{\nabla} \frac{\epsilon_0}{\epsilon(\vec{r})} \end{aligned}$$

↳ charge density with dielectric screening.

↳ surface charge density

$$\begin{aligned} \vec{\nabla} \frac{\epsilon_0}{\epsilon(\vec{r})} &= \hat{z} \frac{\partial}{\partial z} \frac{\epsilon_0}{\epsilon(z)} = \hat{z} \frac{\partial}{\partial z} \begin{cases} \frac{\epsilon_0}{\epsilon_2}, & z < 0 \\ \frac{\epsilon_0}{\epsilon_1}, & z > 0 \end{cases} \\ &= \hat{z} \frac{\partial}{\partial z} \left[\frac{\epsilon_0}{\epsilon_2} \theta(-z) + \frac{\epsilon_0}{\epsilon_1} \theta(z) \right] \\ &= \hat{z} \epsilon_0 \left(\frac{1}{\epsilon_1} - \frac{1}{\epsilon_2} \right) \delta(z) \\ &= \hat{z} \frac{\epsilon_0}{\epsilon_1} \frac{\epsilon_2 - \epsilon_1}{\epsilon_2} \delta(z) \end{aligned}$$

$\epsilon_0 = \epsilon_1$ for our situation.

⑫ Using ⑪ in ⑩ we have

$$\begin{aligned} \epsilon(\vec{r}) \vec{E} \cdot \vec{\nabla} \frac{\phi}{\epsilon(\vec{r})} &= \epsilon(z) E_z(x, y, z) \frac{\epsilon_2 - \epsilon_1}{\epsilon_2} \delta(z) \\ &= \frac{\epsilon(+\delta)}{\epsilon_1} E_z(x, y, +\delta) \frac{\epsilon_2 - \epsilon_1}{\epsilon_2} \delta(z) \end{aligned}$$

and we identify the surface charge density,

$$\sigma(x, y) = \frac{\epsilon_1}{\epsilon_2} (\epsilon_2 - \epsilon_1) E_z(x, y, +\delta)$$

Note $D_z = \epsilon E_z$ is continuous across the slab. Thus, we can use any side to evaluate it

⑬ Using the surface charge density of ⑫ in ⑨ we have the force on the dielectric

$$F = \int_{-\infty}^{+\infty} dx \int_{-\infty}^{+\infty} dy \left[\frac{\epsilon_1}{\epsilon_2} (\epsilon_2 - \epsilon_1) E_z(x, y, +\delta) \right] E_z(x, y, 0)$$

⑭ The electric field, z-component, at $z=0$, is interpreted as.

$$\begin{aligned} E_z(x, y, 0) &= \frac{1}{2} \left[E_z(x, y, +\delta) + E_z(x, y, -\delta) \right] \\ &= \frac{1}{2} \left[E_z(x, y, +\delta) + \frac{\epsilon_1}{\epsilon_2} E_z(x, y, +\delta) \right] \\ &= \frac{\epsilon_1 + \epsilon_2}{2\epsilon_2} E_z(x, y, +\delta) \end{aligned}$$

$$\epsilon_2 E(-\delta) = \epsilon_1 E(+\delta)$$

(15) Using (14) in (13) we have.

$$F = \int_{-\infty}^{+\infty} dx \int_{-\infty}^{+\infty} dy \frac{\epsilon_1}{\epsilon_2} (\epsilon_2 - \epsilon_1) \frac{\epsilon_1 + \epsilon_2}{2\epsilon_2} E_z(x, y, +d)^2$$

(16) Using eqn. (6) in lecture dated 2014 Nov 14 we have

$$E_z(x, y, +d) = \frac{q}{4\pi\epsilon_1} \frac{(-d) \left[1 - \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2} \right]}{(x^2 + y^2 + d^2)^{\frac{3}{2}}} = -\frac{q}{4\pi\epsilon_1} \frac{2\epsilon_2}{\epsilon_1 + \epsilon_2} \frac{d}{(x^2 + y^2 + d^2)^{\frac{3}{2}}}$$

$$\begin{aligned} (17) \quad F &= \int_{-\infty}^{+\infty} dx \int_{-\infty}^{+\infty} dy \frac{\epsilon_1}{\epsilon_2} (\epsilon_2 - \epsilon_1) \frac{\epsilon_1 + \epsilon_2}{2\epsilon_2} \left[-\frac{q}{4\pi\epsilon_1} \frac{2\epsilon_2}{\epsilon_1 + \epsilon_2} \frac{d}{(x^2 + y^2 + d^2)^{\frac{3}{2}}} \right]^2 \\ &= \frac{1}{4\pi\epsilon_1} \frac{\epsilon_2 - \epsilon_1}{\epsilon_1 + \epsilon_2} \frac{q^2}{2\pi} \int_0^\infty r dr \int_0^{2\pi} d\phi \frac{d^2}{(r^2 + d^2)^3} \\ &= \frac{1}{4\pi\epsilon_1} \frac{\epsilon_2 - \epsilon_1}{\epsilon_1 + \epsilon_2} \frac{q^2}{d^2} \int_0^\infty \frac{x dx}{(1 + x^2)^3} \quad \int_0^\infty \frac{x dx}{(1 + x^2)^3} \quad \begin{matrix} 1+x^2 = t \\ 2x dx = dt \end{matrix} \\ &= \frac{1}{4\pi\epsilon_1} \frac{\epsilon_2 - \epsilon_1}{\epsilon_1 + \epsilon_2} \frac{q^2}{d^2} \int_1^\infty \frac{\frac{dt}{2}}{t^3} \\ &= \frac{1}{4\pi\epsilon_1} \frac{\epsilon_2 - \epsilon_1}{\epsilon_1 + \epsilon_2} \frac{q^2}{d^2} \left[-\frac{1}{2t^2} \right]_1^\infty = \frac{1}{4} \\ &= \frac{1}{4\pi\epsilon_1} \frac{\epsilon_2 - \epsilon_1}{\epsilon_1 + \epsilon_2} \frac{q^2}{(2d)^2} \end{aligned}$$

which verifies the expression we calculated in (8). The charge in sign is because the force in (8) is on the charge and the force above is on the dielectric material. They satisfy Newton's third law.

(18) Homework: Calculate the total induced charge to be

$$\int_{-\infty}^{+\infty} dx \int_{-\infty}^{+\infty} dy \sigma(x, y) = \frac{\epsilon_1 - \epsilon_2}{\epsilon_1 + \epsilon_2} q.$$