

① We have defined Bessel's function using an integral representation,

$$J_m(t) = \int_0^{2\pi} \frac{d\alpha}{2\pi} e^{it \cos \alpha} e^{-im\alpha}.$$

It has the series representation

$$\begin{aligned} J_m(t) &= \sum_{n=0}^{\infty} \frac{1}{n!} \frac{(-1)^n}{(n+m)!} \left(\frac{t}{2}\right)^{m+2n} \\ &= \frac{1}{m!} \left(\frac{t}{2}\right)^m - \frac{1}{(m+1)!} \left(\frac{t}{2}\right)^{m+2} + \dots \end{aligned}$$

② The free Green's function

$$G_0(\vec{r}, \vec{r}') = \frac{1}{4\pi\epsilon_0} \sum_{m=-\infty}^{+\infty} e^{im(\phi-\phi')} \int_0^{\infty} dk_z J_m(k_z \rho) J_m(k_z \rho') e^{-k_z |z-z'|}$$

satisfies the Laplacian

$$-\nabla^2 G_0(\vec{r}, \vec{r}') = \delta^{(3)}(\vec{r} - \vec{r}'),$$

which for  $\vec{r}' = 0$  and  $\vec{r} \neq \vec{r}'$  is

$$\nabla^2 G_0(\vec{r}, 0) = 0.$$

Thus, we have.

$$\nabla^2 \left[ J_m(k_z \rho) e^{im\phi} e^{-k_z |z|} \right] = 0$$

③ Using

$$\nabla^2 = \frac{1}{s} \frac{\partial}{\partial s} s \frac{\partial}{\partial s} + \frac{1}{s^2} \frac{\partial^2}{\partial \phi^2} + \frac{\partial^2}{\partial z^2}$$

we have

$$\left[ \frac{1}{s} \frac{\partial}{\partial s} s \frac{\partial}{\partial s} - \frac{m^2}{s^2} + k_z^2 \right] J_m(k_z s) = 0$$

which reads, using  $k_z s = t$ ,

$$\left[ \frac{1}{t} \frac{d}{dt} t \frac{d}{dt} - \frac{m^2}{t^2} + 1 \right] J_m(t) = 0$$

④ In preparation to studying the asymptotic limits of Bessel's function we note,

$$\frac{1}{\sqrt{t}} \frac{d}{dt} \sqrt{t} = \left( \frac{1}{2t} + \frac{d}{dt} \right)$$

$$\left( \frac{1}{\sqrt{t}} \frac{d}{dt} \sqrt{t} \right)^2 = \frac{1}{\sqrt{t}} \frac{d^2}{dt^2} \sqrt{t}$$

$$= \left( \frac{1}{2t} + \frac{d}{dt} \right)^2$$

$$= \frac{1}{4t^2} + \frac{1}{2t} \frac{d}{dt} + \frac{d}{dt} \frac{1}{2t} + \frac{d^2}{dt^2}$$

$$= \frac{1}{4t^2} + \frac{1}{2t} \frac{d}{dt} - \frac{1}{2t^2} + \frac{1}{2t} \frac{d}{dt} + \frac{d^2}{dt^2}$$

$$= \frac{d^2}{dt^2} + \frac{1}{t} \frac{d}{dt} - \frac{1}{4t^2}$$

in terms of which we can write

$$\left[ \frac{1}{\sqrt{t}} \frac{d^2}{dt^2} \sqrt{t} - \frac{(m^2 - \frac{1}{4})}{t^2} + 1 \right] J_m(t) = 0.$$

⑤ We further write

$$\left[ \frac{d^2}{dt^2} - \frac{(m^2 - \frac{1}{4})}{t^2} + 1 \right] \sqrt{t} J_m(t) = 0.$$

⑥ For  $t \rightarrow 0$ , or  $|m^2 - \frac{1}{4}| \gg t^2$ , we have.

$$\frac{d^2}{dt^2} \sqrt{t} J_m(t) = \frac{(m^2 - \frac{1}{4})}{t^2} \sqrt{t} J_m(t).$$

Observing that, for  $m > 0$ ,

$$\frac{d^2}{dt^2} t^{m+\frac{1}{2}} = \frac{d}{dt} (m+\frac{1}{2}) t^{m-\frac{1}{2}} = (m^2 - \frac{1}{4}) t^{m-\frac{3}{2}} = \frac{(m^2 - \frac{1}{4})}{t^2} t^{m+\frac{1}{2}},$$

we find

$$\lim_{t \rightarrow 0} J_m(t) \approx t^m,$$

which is consistent with the more accurate form,

using ①,

$$\lim_{t \rightarrow 0} J_m(t) = \frac{1}{m!} \left(\frac{t}{2}\right)^m.$$

⑦ For  $t \rightarrow \infty$ , or  $|m^2 - \frac{1}{4}| \ll t^2$ , we have.

$$\left( \frac{d^2}{dt^2} + 1 \right) \sqrt{t} J_m(t) = 0$$

which implies  $\sqrt{t} J_m(t)$  is of the form  $\cos t$  and  $\sin t$ . This is true, and the more accurate

form of this limit is

$$\lim_{t \rightarrow \infty} J_m(t) \approx \sqrt{\frac{\pi}{2t}} \cos\left(t - (m + \frac{1}{2})\frac{\pi}{2}\right).$$

⑧ Finally, we present the completeness relation for the Bessel function. We begin by noting

$$\int \frac{d^2 k_\perp}{(2\pi)^2} e^{i \vec{k}_\perp \cdot (\vec{r} - \vec{r}')_\perp} = \frac{1}{2\pi} \sum_{m=-\infty}^{+\infty} e^{im(\phi - \phi')} \int_0^\infty k_\perp dk_\perp J_m(k_\perp r) J_m(k_\perp r'),$$

which is also the delta function in two dimensions,

$$\int \frac{d^2 k_\perp}{(2\pi)^2} e^{i \vec{k}_\perp \cdot (\vec{r} - \vec{r}')_\perp} = \delta^{(2)}(\vec{r}_\perp - \vec{r}'_\perp) = \frac{\delta(r - r')}{r} \delta(\phi - \phi').$$

Then, using

$$\delta(\phi - \phi') = \frac{1}{2\pi} \sum_{m=-\infty}^{+\infty} e^{im(\phi - \phi')}$$

we obtain

$$\frac{\delta(r - r')}{r} = \int_0^\infty k_\perp dk_\perp J_m(k_\perp r) J_m(k_\perp r'),$$

which is the completeness relation for Bessel functions.

⑨ The orthogonality of Bessel functions is given by

$$\frac{\delta(k_\perp - k'_\perp)}{k_\perp} = \int_0^\infty r dr J_m(k_\perp r) J_m(k'_\perp r).$$