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①

$$\textcircled{1} \quad x^\mu = (ct, \vec{x}) \rightarrow dx^\mu dx_\mu = -c^2 dt^2 + d\vec{x} \cdot d\vec{x} = -ds^2$$

$$u^\alpha = c \frac{dx^\alpha}{ds} \\ = (c\gamma, \vec{v}\gamma) \rightarrow u^\alpha u_\alpha = -c^2$$

Physica

$$\left(\frac{E}{c}, \vec{P}\right) \leftarrow P^\alpha = m u^\alpha = (m c \gamma, m \vec{v} \gamma)$$

②

$$a^\alpha = c \frac{du^\alpha}{ds}$$

$$\frac{cdt}{ds} = \gamma$$

$$= c \frac{dt}{ds} \frac{d}{dt} u^\alpha$$

$$= \gamma \frac{d}{dt} (c\gamma, \vec{v}\gamma)$$

$$= \gamma \left(c \frac{d\gamma}{dt}, \frac{d\vec{v}}{dt} \gamma + \vec{v} \frac{d\gamma}{dt} \right)$$

$$= \gamma \left(c \frac{d\gamma}{dt}, \vec{a} \gamma + \vec{v} \frac{d\gamma}{dt} \right)$$

$$\vec{a} = \frac{d\vec{v}}{dt}$$

③ Observe that starting from

$$u^\alpha u_\alpha = -c^2$$

we can show that (by differentiating with respect to s)

$$a^\alpha u_\alpha = 0.$$

HW: Is a^α a space-like or time-like quantity?

④ Thus, using

$$u^\alpha = r(c, \vec{v})$$

$$a^\alpha = r\left(c \frac{dr}{dt}, \vec{a} r + \vec{v} \frac{dr}{dt}\right)$$

we have

$$a^\alpha u_\alpha = 0$$

$$r^2 \left[-c^2 \frac{dr}{dt} + \vec{v} \cdot \vec{a} r + \vec{v} \cdot \vec{v} \frac{dr}{dt} \right] = 0$$

$$-c^2 r^2 \left[\left(1 - \frac{v^2}{c^2}\right) \frac{dr}{dt} - \frac{\vec{v} \cdot \vec{a}}{c^2} r \right] = 0$$

$$1 - \frac{v^2}{c^2} = \frac{1}{r^2}$$

$$-c^2 \left[\frac{dr}{dt} - \frac{\vec{v} \cdot \vec{a}}{c^2} r^3 \right] = 0$$

$$\Rightarrow \frac{dr}{dt} = \frac{\vec{v} \cdot \vec{a}}{c^2} r^3$$

⑤ Using ④ in the expression of a^α in ④ we have

$$a^\alpha = r\left(c \frac{dr}{dt}, \vec{a} r + \vec{v} \frac{dr}{dt}\right)$$

$$= r\left(\frac{\vec{v} \cdot \vec{a}}{c} r^3, \vec{a} r + \vec{v} \frac{\vec{v} \cdot \vec{a}}{c^2} r^3\right)$$

⑥ Let us next define 4-force as

$$\begin{aligned}
 f^\alpha &= c \frac{d}{ds} P^\alpha \\
 &= c \frac{dt}{ds} \frac{d}{dt} P^\alpha \\
 &= \gamma \frac{d}{dt} \left(\frac{E}{c}, \vec{P} \right) \\
 &= \gamma \left(\frac{1}{c} \frac{dE}{dt}, \frac{d\vec{P}}{dt} \right)
 \end{aligned}$$

⑦ Relativistic equation of motion is given using

$$\begin{aligned}
 f^\alpha &= m a^\alpha \\
 \gamma \left(\frac{1}{c} \frac{dE}{dt}, \frac{d\vec{P}}{dt} \right) &= m \gamma \left(c \frac{dr}{dt}, \vec{a} r + \vec{v} \frac{dr}{dt} \right) \\
 &= m \gamma \left(\frac{\vec{v} \cdot \vec{a}}{c} r^3, \vec{a} r + \vec{v} \frac{\vec{v} \cdot \vec{a}}{c^2} r^3 \right)
 \end{aligned}$$

This implies

$$\begin{aligned}
 \frac{dE}{dt} &= \frac{d}{dt} (m c r) = m \frac{\vec{v} \cdot \vec{a}}{c} r^3 \\
 \frac{d\vec{P}}{dt} &= \frac{d}{dt} (m \vec{v} r) = m \vec{a} r + m \vec{v} \frac{\vec{v} \cdot \vec{a}}{c^2} r^3
 \end{aligned}$$

⑧ Let us calculate $a^\alpha a_\alpha$, which will remain invariant between observers.

$$\begin{aligned}
 a^\alpha &= r \left(\frac{\vec{v} \cdot \vec{a}}{c} r^3, \vec{a} r + \vec{v} \frac{\vec{v} \cdot \vec{a}}{c^2} r^3 \right) \\
 a^\alpha a_\alpha &= r^2 \left[- \left(\frac{\vec{v} \cdot \vec{a}}{c} \right)^2 r^6 + \left(\vec{a} r + \vec{v} \frac{\vec{v} \cdot \vec{a}}{c^2} r^3 \right)^2 \right] \\
 &= r^2 \left[- \left(\frac{\vec{v} \cdot \vec{a}}{c} \right)^2 r^6 + \vec{a} \cdot \vec{a} r^2 + \frac{v^2}{c^2} \left(\frac{\vec{v} \cdot \vec{a}}{c} \right)^2 r^6 + 2 \left(\frac{\vec{v} \cdot \vec{a}}{c} \right)^2 r^4 \right] \\
 &= r^2 \left[\vec{a} \cdot \vec{a} r^2 + \left(\frac{\vec{v} \cdot \vec{a}}{c} \right)^2 r^4 \left\{ 2 + \underbrace{\beta^2 r^2 - r^2}_{-1} \right\} \right] \\
 &= \vec{a} \cdot \vec{a} r^4 + \left(\frac{\vec{v} \cdot \vec{a}}{c} \right)^2 r^6
 \end{aligned}$$

⑨ Using $(\vec{v} \times \vec{a}) \cdot (\vec{v} \times \vec{a}) = \vec{v} \cdot \vec{v} \vec{a} \cdot \vec{a} - (\vec{v} \cdot \vec{a})^2$

we can also write

$$\begin{aligned}
 a^\alpha a_\alpha &= \vec{a} \cdot \vec{a} r^4 + \left[\frac{v^2}{c^2} \vec{a} \cdot \vec{a} - \left(\frac{\vec{v} \times \vec{a}}{c} \right) \cdot \left(\frac{\vec{v} \times \vec{a}}{c} \right) \right] r^6 \\
 &= \vec{a} \cdot \vec{a} r^4 \underbrace{(1 - \beta^2 r^2)}_{r^2} - \left(\frac{\vec{v} \times \vec{a}}{c} \right) \cdot \left(\frac{\vec{v} \times \vec{a}}{c} \right) r^6 \\
 &= \vec{a} \cdot \vec{a} r^6 - \left(\frac{\vec{v} \times \vec{a}}{c} \right) \cdot \left(\frac{\vec{v} \times \vec{a}}{c} \right) r^6
 \end{aligned}$$

(10) Using (8) and (9) we conclude.

$$\vec{v} \perp \vec{a} : \quad a^\alpha a_\alpha = \vec{a} \cdot \vec{a} r^4$$

$$\vec{v} \parallel \vec{a} : \quad a^\alpha a_\alpha = \vec{a} \cdot \vec{a} r^6$$

Motion in a magnetic field (under simple initial condition) corresponds to the case $\vec{v} \perp \vec{a}$, and motion in an electric field (under simple initial condition) corresponds to the case $\vec{v} \parallel \vec{a}$.