

## Homework No. 09 (2024 Spring)

### PHYS 510: CLASSICAL MECHANICS

*School of Physics and Applied Physics, Southern Illinois University–Carbondale*

Due date: Tuesday, 2024 Apr 17, 4.30pm

1. (20 points.) For two functions

$$A = A(\mathbf{x}, \mathbf{p}, t), \quad (1a)$$

$$B = B(\mathbf{x}, \mathbf{p}, t), \quad (1b)$$

the Poisson bracket with respect to the canonical variables  $\mathbf{x}$  and  $\mathbf{p}$  is defined as

$$[A, B]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} \equiv \frac{\partial A}{\partial \mathbf{x}} \cdot \frac{\partial B}{\partial \mathbf{p}} - \frac{\partial A}{\partial \mathbf{p}} \cdot \frac{\partial B}{\partial \mathbf{x}}. \quad (2)$$

Show that the Poisson bracket satisfies the conditions for a Lie algebra. That is, show that

- (a) Antisymmetry:

$$[A, B]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} = -[B, A]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}}. \quad (3)$$

- (b) Bilinearity: ( $a$  and  $b$  are numbers.)

$$[aA + bB, C]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} = a[A, C]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} + b[B, C]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}}. \quad (4)$$

Further show that

$$[AB, C]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} = A[B, C]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} + [A, C]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} B. \quad (5)$$

- (c) Jacobi's identity:

$$[A, [B, C]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}}]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} + [B, [C, A]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}}]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} + [C, [A, B]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}}]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} = 0. \quad (6)$$

2. (20 points.) Show that the commutator of two matrices,

$$[\mathbf{A}, \mathbf{B}] \equiv \mathbf{AB} - \mathbf{BA}, \quad (7)$$

satisfies the conditions for a Lie algebra, as does the Poisson bracket. In particular show that

- (a) Antisymmetry:

$$[\mathbf{A}, \mathbf{B}] = -[\mathbf{B}, \mathbf{A}]. \quad (8)$$

(b) Bilinearity: ( $a$  and  $b$  are numbers.)

$$[a\mathbf{A} + b\mathbf{B}, \mathbf{C}] = a[\mathbf{A}, \mathbf{C}] + b[\mathbf{B}, \mathbf{C}]. \quad (9)$$

Further show that

$$[\mathbf{AB}, \mathbf{C}] = \mathbf{A}[\mathbf{B}, \mathbf{C}] + [\mathbf{A}, \mathbf{C}]\mathbf{B}. \quad (10)$$

(c) Jacobi's identity:

$$[\mathbf{A}, [\mathbf{B}, \mathbf{C}]] + [\mathbf{B}, [\mathbf{C}, \mathbf{A}]] + [\mathbf{C}, [\mathbf{A}, \mathbf{B}]] = 0. \quad (11)$$

3. (20 points.) Show that the vector product of two vectors, in this problem denoted using

$$[\mathbf{A}, \mathbf{B}]_v \equiv \mathbf{A} \times \mathbf{B}, \quad (12)$$

satisfies the conditions for a Lie algebra, as does the Poisson bracket. In particular show that

(a) Antisymmetry:

$$[\mathbf{A}, \mathbf{B}]_v = -[\mathbf{B}, \mathbf{A}]_v. \quad (13)$$

(b) Bilinearity: ( $a$  and  $b$  are numbers.)

$$[a\mathbf{A} + b\mathbf{B}, \mathbf{C}]_v = a[\mathbf{A}, \mathbf{C}]_v + b[\mathbf{B}, \mathbf{C}]_v. \quad (14)$$

Further show that

$$[\mathbf{A} \times \mathbf{B}, \mathbf{C}]_v = \mathbf{A} \times [\mathbf{B}, \mathbf{C}]_v + [\mathbf{A}, \mathbf{C}]_v \times \mathbf{B}. \quad (15)$$

(c) Jacobi's identity:

$$[\mathbf{A}, [\mathbf{B}, \mathbf{C}]_v]_v + [\mathbf{B}, [\mathbf{C}, \mathbf{A}]_v]_v + [\mathbf{C}, [\mathbf{A}, \mathbf{B}]_v]_v = 0. \quad (16)$$

4. (20 points.) Given  $F$  and  $G$  are constants of motion, that is

$$[F, H]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} = 0 \quad \text{and} \quad [G, H]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} = 0. \quad (17)$$

Then, using Jacobi's identity, show that  $[F, G]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}}$  is also a constant of motion. Thus, conclude the following:

(a) If  $L_x$  and  $L_y$  are constants of motion, then  $L_z$  is also a constant of motion.

(b) If  $p_x$  and  $L_z$  are constants of motion, then  $p_y$  is also a constant of motion.

5. (20 points.) (Refer Sec. 21 Dirac's QM book.)

The product rule for Poisson bracket can be stated in the following different forms:

$$[A_1 A_2, B]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} = A_1 [A_2, B]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} + [A_1, B]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} A_2, \quad (18a)$$

$$[A, B_1 B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} = B_1 [A, B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} + [A, B_1]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} B_2. \quad (18b)$$

(a) Thus, evaluate, in two different ways,

$$\begin{aligned} [A_1 A_2, B_1 B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} &= A_1 B_1 [A_2, B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} + A_1 [A_2, B_1]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} B_2 \\ &\quad + B_1 [A_1, B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} A_2 + [A_1, B_1]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} B_2 A_2, \end{aligned} \quad (19a)$$

$$\begin{aligned} [A_1 A_2, B_1 B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} &= B_1 A_1 [A_2, B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} + B_1 [A_1, B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} A_2 \\ &\quad + A_1 [A_2, B_1]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} B_2 + [A_1, B_1]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} A_2 B_2. \end{aligned} \quad (19b)$$

(b) Subtracting these results, obtain

$$(A_1 B_1 - B_1 A_1) [A_2, B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} = [A_1, B_1]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} (A_2 B_2 - B_2 A_2). \quad (20)$$

Thus, using the definition of the commutation relation,

$$[A, B] \equiv AB - BA, \quad (21)$$

obtain the relation

$$[A_1, B_1] [A_2, B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} = [A_1, B_1]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}} [A_2, B_2]. \quad (22)$$

(c) Since this condition holds for  $A_1$  and  $B_1$  independent of  $A_2$  and  $B_2$ , conclude that

$$[A_1, B_1] = i\hbar [A_1, B_1]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}}, \quad (23a)$$

$$[A_2, B_2] = i\hbar [A_2, B_2]_{\mathbf{x}, \mathbf{p}}^{\text{P.B.}}, \quad (23b)$$

where  $i\hbar$  is necessarily a constant, independent of  $A_1$ ,  $A_2$ ,  $B_1$ , and  $B_2$ . This is the connection between the commutator bracket in quantum mechanics and the Poisson bracket in classical mechanics. If  $A$ 's and  $B$ 's are numbers, then, because their commutation relation is equal to zero, we necessarily have  $\hbar = 0$ . But, if the commutation relation of  $A$ 's and  $B$ 's is not zero, then finite values of  $\hbar$  is allowed.

(d) Here the imaginary number  $i = \sqrt{-1}$ . Show that the constant  $\hbar$  is a real number if we presume the Poisson bracket to be real, and require the construction

$$C = \frac{1}{i}(AB - BA) \quad (24)$$

to be Hermitian. Experiment dictates that  $\hbar = h/2\pi$ , where

$$h \sim 6.63 \times 10^{-34} \text{ J}\cdot\text{s} \quad (25)$$

is the Planck's constant with dimensions of action.